

Nambu bracket and sigma-models

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Chern-Simons model

The Poisson sigma-model and its twisted version

The Nambu mechanics

The Nambu σ -models

Put M -3d manifold with a principal bundle with group G (A -conection form)

$$S_{CS} = \frac{k}{4\pi} \int_M \text{tr}(AdA + \frac{2}{3}A \wedge A \wedge A)$$

The partition function

$$Z(k) = \int DA e^{iS_{CS}}$$

Feynman's integral over all equivalence classes of connections modulo gauge transformations

k is a coupling constant and for gauge invariance $k \in \mathbb{Z}$

Here gauge invariant observables are the Wilson lines

$$W_R(C) = \text{tr}_R P e^{\oint_C A},$$

R is irreducible representation of G , C is oriented closed curve (knot), P is P-ordering

Put a link C with components K_1, \dots, K_L

$$W_{R_1, \dots, R_L}(C) = \langle W_{R_1}^{K_1} \dots W_{R_L}^{K_L} \rangle = \frac{1}{Z(M)} \int DA \left(\prod_{i=1}^L W_{R_i}^{K_i} \right) e^{iS_{CS}}$$

For $SU(N)$ one can obtain the HOMFLY-PT polynomial $P_C(q, \lambda)$ (for $SU(2)$ the Jones polynomial) and all $R_i = \square$

$$W_{\square, \dots, \square}(C) \sim P_C(q, \lambda),$$

where $q = \exp(\frac{2\pi i}{k+N})$ and $\lambda = q^N$

$SU(2) \rightarrow$ the Jones polynomials

$SO(N) \rightarrow$ the Kauffman polynomial

The partition functions for $SU(N)$

The partition function for $SU(N)$

$$Z(S^3, N, k) = \frac{1}{\sqrt{N}(k+N)^{\frac{N}{2}-1}} \prod_{i=1}^{N-1} \left(2 \sin\left(\frac{\pi i}{k+N}\right)\right)^{N-i}$$

$Z(T^3, N, k)$ counts the number of integrable irreducible highest-weight representations of the affine Kac-Moody algebra $SU(N)_k$

$$Z(T^3, N, k) = \frac{1}{B(N, K)}$$

The Courant sigma-model (D. Roytenberg '01)

Put some vector bundle $E \rightarrow M_0$. Let us consider supermanifold ΠT^*E . If $\{x_i\}$ are local coordinates on M_0 and $\{e_i\}$ is a local basis of section for E with $\langle e_a, e_b \rangle = g_{ab} = \text{const}$.

Courant algebroid $(E, \langle \cdot, \cdot \rangle)$ the anchor map has matrix P_a^i

$$S(X, A, F) = \int_N F_i dX^i + \frac{1}{2} A^a g_{ab} dA^b - A^a P_a^i(X) F_i + \frac{1}{6} T_{abc} A^a A^b A^c + \dots$$

$F = 0$ and $T_{abc} = \epsilon_{abc}$ one can obtain usual Chern-Simons theory

The Poisson manifold

Put N is a Poisson manifold

(N has Poisson bi-vector $\pi = \frac{1}{2}\pi^{ij}(x)\partial_i \wedge \partial_j$)

$$\{f(x), g(x)\} = \pi^{ij}(x) \frac{\partial f}{\partial x_i} \frac{\partial g}{\partial x_j}$$

The Jacobi identity $\{\{f, g\}, h\} + cycl. = 0$

$$[\pi, \pi]_{NS} = 0$$

The time evolution is given by

$$\frac{df}{dt} = \{H, f\}$$

The Schouten-Nijenhuis bracket

The bracket for polivectors

$$[\cdot, \cdot]_{NS} : M^n(N) \times M^m(N) \rightarrow M^{n+m-1}(N),$$

$$[u_1 \wedge \dots \wedge u_n, v_1 \wedge \dots \wedge v_m]_{NS} = \sum_{i,j} (-1)^{i+j} [u_i, v_j] \wedge \dots \wedge \cancel{u_i} \wedge \dots \wedge v_1 \wedge \dots \wedge \cancel{v_j} \wedge \dots \wedge v_n$$

Usually 2d sigma-models have action

$$S = \int_{\Sigma} \eta^{\mu\nu}(x) g^{ij}(\phi) \partial_{\mu} \phi_i \partial_{\nu} \phi_j \sqrt{\eta(x)} d^2x$$

here Σ is two-dimensional (super)manifold (world-sheet or space-time) with metric η and fields ϕ_i lie in a (super)manifold L (target space) with metric g

Simple example $O(N)$ -model

$$S = \frac{1}{2g^2} \int_{\mathcal{X}} \partial_{\mu} \phi_i \partial_{\mu} \phi_i, \quad \phi_i \phi_i = 1$$

Supersymmetric examples are A and B-models

The Poisson σ -model

Put (L, π) the Poisson manifold with fields 0-forms X^i (coordinates on target L) and 1-form A

$$A = A_{\mu i} dx^\mu \wedge dX^i$$

Then the topological Poisson σ -model

$$S_t = \int_{\Sigma} A_i \wedge dX^i + \frac{1}{2} \pi^{ij}(X) A_i \wedge dA_j$$

$$S_1 = \int_{\Sigma} A_i \wedge dX^i + \frac{1}{2} \pi^{ij} A_i \wedge A_j + \frac{1}{2} (G^{-1})^{ij} A_i \wedge *A_j$$

Put $2\pi\alpha' = 1$. The string sigma-model is

$$S_2 = \frac{1}{2} \int_{\Sigma} (g_{ij} dX^i \wedge *dX^j + B_{ij} dX^i \wedge dX^j)$$

g and B are related to G^{-1} and π

$$\frac{1}{g + B} = G^{-1} + \pi$$

Topological limit $G^{-1} \rightarrow 0$

Noncommutative limit $g \rightarrow 0 (\alpha' \rightarrow 0)$

$$x_i \rightarrow \tilde{x}_i = x_i + \theta_{ij} A_j$$

If (N, π) has a closed 3-form H . Then the twisted Jacobi identity

$$\frac{1}{2}[\pi, \pi]_{SN} = \Lambda^3 \pi H$$

$$S(X, A) = \int_{\Sigma} (A_i \wedge dX^i + \frac{1}{2} \pi^{ij} A_i \wedge A_j) + S_{WZ},$$

$$S_{WZ} = \int_{\Lambda} \tilde{H},$$

where $\partial\Lambda = \Sigma$ and \tilde{H} mean the pullback of H to Λ by extension to V of the map X^i

an algebra with n-ary map $[] : A^{\wedge n} \rightarrow A$ (Nambu-Filippov algebra)

$$[L_1, L_2, \dots, L_n] = (-1)^{\epsilon(\sigma)} [L_{\sigma_1}, \dots, L_{\sigma_n}] \quad \sigma \in S_n$$

$$[L_{i_1}, L_{i_2}, \dots, L_{i_n}] = f_{i_1 \dots i_n}^k L_k$$

The Leibniz rule

$$[L_0 L_{i_1}, L_{i_2}, \dots, L_{i_n}] = L_0 [L_{i_1}, L_{i_2}, \dots, L_{i_n}] + [L_0, L_{i_2}, \dots, L_{i_n}] L_{i_1}$$

There is generalized Jacobi identity (Fundamental identity)

Example. Nambu-Heisenberg algebra

$$[x_1, x_2, x_3] = x_1 x_2 x_3 - x_1 x_3 x_2 + x_3 x_1 x_2 - x_3 x_2 x_1 + x_2 x_3 x_1 - x_2 x_1 x_3 = -i\hbar I$$

The Nambu bracket

$$\{f_1, f_2, \dots, f_n\} = \pi^n(df_1, df_2, \dots, df_n),$$

$$\pi^n = \pi^{i_1 \dots i_n} \partial_{i_1} \wedge \dots \wedge \partial_{i_n}$$

Generalized Jacobi identity (Fundamental identity)

$$\begin{aligned} & \{ \{f_1, f_2, \dots, f_n\}, f_{n+1}, \dots, f_{2n-1} \} + \{ f_n, \{f_1, f_2, \dots, f_n\}, f_{n+2}, \dots, f_{2n-1} \} + \dots \\ & + \{ f_{n+1}, \dots, f_{2n-1}, \{f_1, f_2, \dots, f_n\} \} = \{ f_1, f_2, \dots, f_{n-1}, \{f_n, \dots, f_{2n-1}\} \} \end{aligned}$$

$$\frac{df}{dt} = \{H_1, H_2, \dots, H_n, f\}$$

Example of Nambu mechanics

The asymmetric Euler top with $H_1 = \frac{1}{2}\left(\frac{L_1^2}{I_1} + \frac{L_2^2}{I_2} + \frac{L_3^2}{I_3}\right)$ and $H_2 = \frac{1}{2}(L_1^2 + L_2^2 + L_3^2)$

The Nambu equations $\frac{dL_i}{dt} = \epsilon^{ijk} \partial_j H_1 \partial_k H_2$

$$\frac{dL_1}{dt} = \left(\frac{1}{I_2} - \frac{1}{I_3}\right) L_2 L_3, \quad \frac{dL_2}{dt} = \left(\frac{1}{I_3} - \frac{1}{I_1}\right) L_3 L_1, \quad \frac{dL_3}{dt} = \left(\frac{1}{I_2} - \frac{1}{I_1}\right) L_1 L_2$$

3d Nambu topological sigma-model (Schupp, Jurco '2012)

Put 1-forms $A_{ij} = A_{ij\alpha} dx^\alpha$ and 2-forms $B_i = B_{i\alpha\beta} dx^\alpha \wedge dx^\beta$ on N and $\alpha, \beta = 0, 1, 2$

The action for the topological Nambu sigma-model has form

$$S_N(X, A, B) = \int_N (B_i \wedge dX^i + A_{ij} \wedge dX^i dX^j + \frac{1}{6} \pi^{ijk} A_{ij} \wedge B_k)$$

Here one can add terms which depends on metric

$$\frac{1}{2} \int_N (G^{ij} B_i \wedge *B_j + \tilde{G}^{klmn} A_{kl} \wedge *A_{mn})$$

There is a similar description by the membrane sigma model and Seiberg-Witten map

Remark about twisted Nambu bracket and sigma-model

It is possible to imagine the Jacobi identity for the Nambu bracket like usual the Poisson bracket

$$[\pi^n, \pi^n]_{SN} = 0$$

If manifold V ($N = \partial V$) has some $(n+1)$ -form H we again can add right side in the Jacobi identity and Wess-Zumino term in the action

$$[\pi^n, \pi^n]_{SN} = \Lambda^3 \pi^n H.$$

$$S = S_N + \int_V \tilde{H}$$

Thank you for attention